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# Geodesic Motion in the Schwarzschild Metric

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## The Geodesic Equations for Massless Particles

The geodesic motion of freely falling particles can be obtained from the solution of the geodesic equations. We choose the coordinates so that the motion is confined in the  $\theta = \frac{\pi}{2}$  plane. Therefore

$\dot{\theta} = 0$ ,  $\sin(\theta) = 1$ . Then we have to solve the system

$$\ddot{r} = \frac{r^2 - e^2}{r^2(1 - \frac{2}{r})} + \frac{1}{r^3} \left(1 - \frac{2}{r}\right)$$

$$\dot{t} = \frac{e}{1 - \frac{2}{r}}$$

$$\dot{\phi} = \frac{1}{r^2}$$

where  $\dot{t} = \frac{d}{d\lambda} t(\lambda)$  etc, and the conserved quantities

$$e = \left(1 - \frac{2}{r}\right) \dot{t}$$

$$1 = r^2 \dot{\phi}$$

are given.

We have rescaled the affine parameter of the null geodesic, so that the angular momentum  $l = 1$ .

The radial motion “conservation of energy” is given by

$$e^2 = \frac{1}{2} \dot{r}^2 + W_{\text{eff}}(r)$$

$$W_{\text{eff}}(r) = \frac{1}{r^2} \left(1 - \frac{2}{r}\right)$$

The turning points of the radial motion are calculated from the equation

$$W_{\text{eff}}(r) = e^2$$

The maximum of  $W_{\text{eff}}(r)$  is at  $r_{\max} = 3$  and  $W_{\text{eff}}(r_{\max}) = \frac{1}{27}$ . Therefore, we have circular orbits when  $e^2 = \frac{1}{27}$ ,  $b = \sqrt{27}$

The impact parameter is

$$b^2 = \frac{1}{e^2}$$

In the above equations, dimensionless quantities are used. To go back to geometrized units, use the

dictionary:

First restore l:

$$\lambda \rightarrow l \lambda \quad \frac{d}{d\lambda} \rightarrow \frac{1}{l} \frac{d}{dl}$$

$$t \rightarrow l^{-1} t \Rightarrow e \rightarrow l^{-1} e$$

Then restore M:

$$\tau \rightarrow \frac{\tau}{M} \quad t \rightarrow \frac{t}{M} \quad r \rightarrow \frac{r}{M} \quad \phi \rightarrow \phi$$

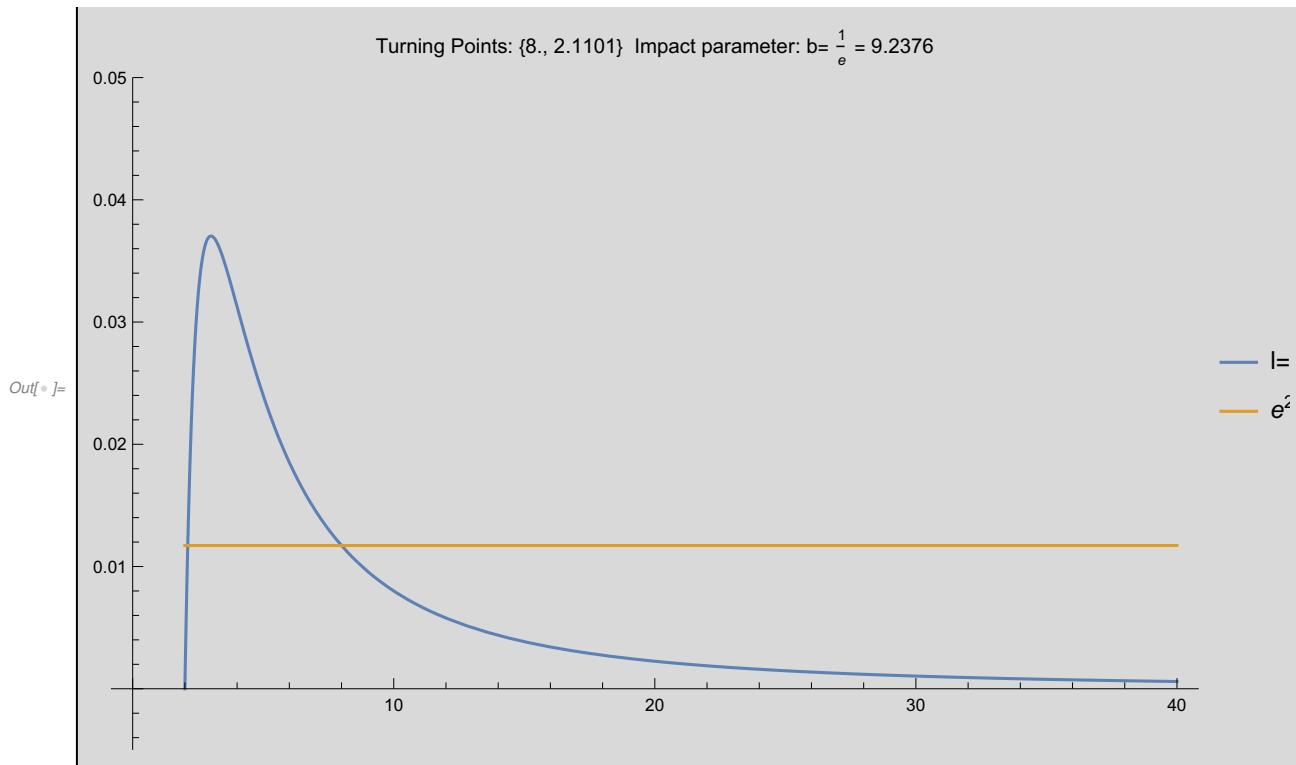
$$l \rightarrow \frac{l}{M} \quad W_{\text{eff}} \rightarrow M^2 W_{\text{eff}} \quad e \rightarrow e \quad b \rightarrow b$$

## Effective potential for radial motion of massless particle

```
In[1]:= l = 4.3; rmax = 8.;

Weff[r_] := 1/(r^2) (1 - 2/r);

Plot[Evaluate[{Weff[r], Weff[rmax]}, {r, 2, 40}], PlotRange -> {-0.005, 0.05}, PlotLegends -> {"l=" <> ToString[l], "e^2= " <> ToString[Weff[rmax]]}, PlotLabel -> "Turning Points: " <> ToString[Select[r /. NSolve[Weff[r] == Weff[rmax], r], # ∈ Reals && # > 0 &]] <> " Impact parameter: b= 1/e = " <> ToString[1/Sqrt[Weff[rmax]]], ImageSize -> Large]
```



## Geodesic Equations

```

lnf :=  $W_{eff}[r] := \frac{1}{r^2} \left(1 - \frac{2}{r}\right);$ 

solveGeodesicEqs[e_, phi0_, r0_, vr0_, lambdaFinal_] := NDSolve[{
    t'[tau] ==  $\frac{e}{1 - \frac{2}{r[\tau]}}$ ,
    phi'[tau] ==  $\frac{1}{r[\tau]^2}$ ,
    r''[tau] ==  $-\frac{e^2}{r[\tau]^2 \left(1 - \frac{2}{r[\tau]}\right)} + \frac{(r'[tau])^2}{r[\tau]^2 \left(1 - \frac{2}{r[\tau]}\right)} + \left(1 - \frac{2}{r[\tau]}\right) \frac{1}{r[\tau]^3}$ ,
    t[0] == 0, phi[0] == phi0, r[0] == r0,
    r'[0] == vr0 (* initial conditions *)
}, {t, phi, r}, {tau, 0, lambdaFinal}]
];

```

## Plotting function

Skip this section if you are not interested in the plotting commands.

This is the function that plots the results.

```
In[1]:= plotResults[xyRange_] :=
(* xyRange is the range of the x-y plot *)
Module[{xyrange = xyRange},
GraphicsGrid[
{
(* Arrange the plots in a grid *)
{
(* 1st row *)
Show[{
ParametricPlot[Evaluate[
{r[\tau] Cos[\phi[\tau]], r[\tau] Sin[\phi[\tau]]} (* x=r cos(phi) y=r sin(phi) *)
/. sol], {\tau, 0.1, \lambda_{max}}, AxesLabel -> {"x", "y"}, PlotRange -> xyrange],
Graphics[{Red, Disk[{0, 0}, 2]}]
(* a red disk in the black hole area *)
}],
Plot[{Weff[r], Energy
(* The effective potential with the energy level E *)
}, {r, 0, 40},
PlotRange -> {-0.001, 0.05},
PlotLabel -> "Turning Points: " (* compute
the turning points *)
<> ToString[Select[r /. NSolve[Weff[r] == Energy, r], # ∈ Reals && # > 0 &]]
<> " e^2= " <> ToString[Energy]
(* display the value of E on the plot *)
<> " b = " <> ToString[ $\frac{1}{e}$ ],
AxesLabel -> {"r", "Weff[r]"}]
}, (* 1st row *)
{(* 2nd row *)
Plot[Evaluate[{
r[\tau]
(* radial coordinate as a function of λ *)
/. sol], {\tau, 0, \lambda_{max}}, AxesLabel -> {"λ", "r[λ]"}, PlotRange -> All],
Plot[Evaluate[{
Mod[\phi[\tau], 2 π]]
```

```

(* angular coordinate as a function of λ *)  

} /. sol], {τ, 0, λmax}, AxesLabel → {"λ", "ϕ[λ]"}, PlotRange → All]  

}, (* 2nd row *)  

{ (* 3rd row *)  

Plot[Evaluate[{  

    t[τ]  

    (* time coordinate as a function of λ *)  

} /. sol], {τ, 0, λmax}, AxesLabel → {"λ", "t[λ]"}],  

Plot[Evaluate[  

    (* check numerical errors in ε *)  

    (r'[τ]^2 + Weff[r[τ]] - Energy) /. sol  

], {τ, 0, λmax}, AxesLabel → {"λ", "Δε"}, PlotLabel → "Error in ε ≈ e^2"]  

} (* 3rd row *)  

}  

}, ImageSize → Full]
]

```

## Solutions to the geodesic equations for massless particles

Set:

r1: one of the turning points  
 e :  $b^2 = 1/e^2$   
 r0: initial radial position  
 λmax: maximum time for numerical integration for the affine parameter λ  
 radialdirection: the sign of the initial radial velocity  $\dot{r}(0)$

r1 will determine the conserved quantity e, which in turn will determine  $|\dot{r}(0)|$   
 $v_0 = (\text{radial direction}) \times |\dot{r}(0)|$

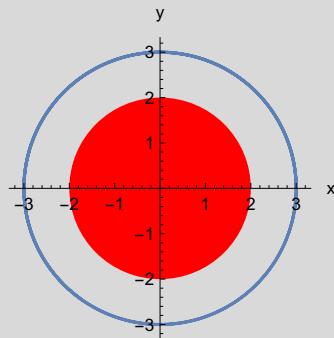
### Circular trajectories

The unstable circular orbit for  $r(0) = r_{\max} = 3$ ,  $e^2 = W_{\text{eff}}(r_{\max}) = \frac{1}{\sqrt{27}}$

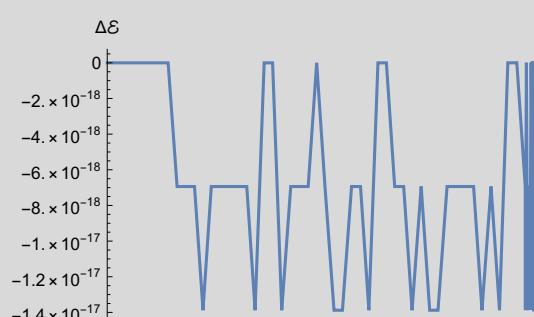
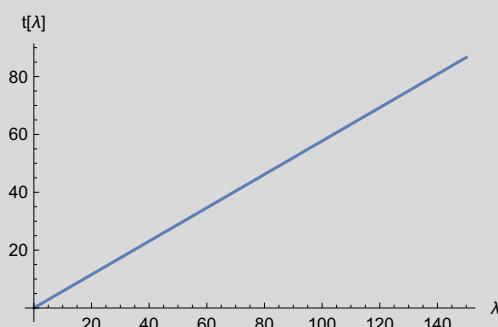
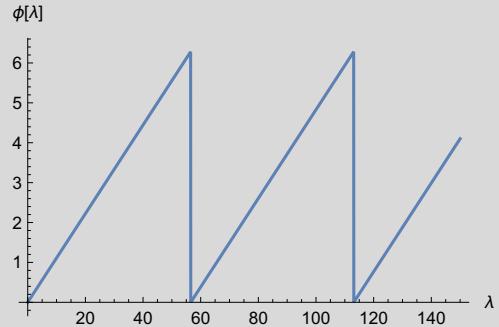
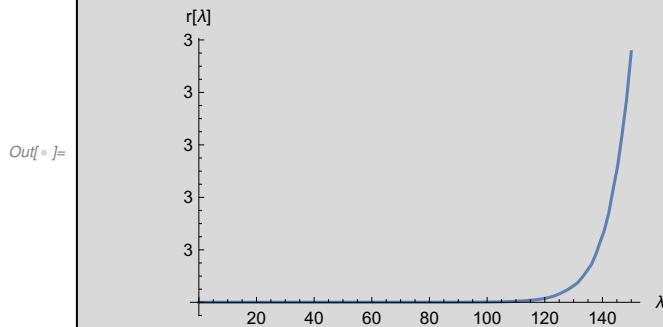
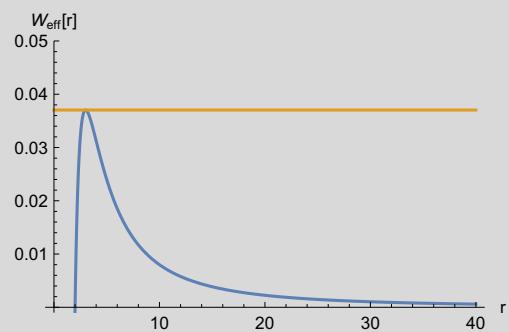
The numerical integration has an instability for  $τ \approx 159$

```
In[6]:= r1 = 3.; λmax = 150;
r0 = 3.; radialdirection = -1; (* Set it to +/- 1. The sign of ḙ(0) *)
Energy = Weff[r1]; e = Sqrt[Energy]; b = 1/e;
v0 = radialdirection Sqrt[Energy - Weff[r0]];
ϕ0 = 0;

sol = solveGeodesicEqs[e, ϕ0, r0, v0, λmax]; gg = plotResults[All]
```



Turning Points: {3., 3.}  $e^2 = 0.037037$   $b = 5.19615$

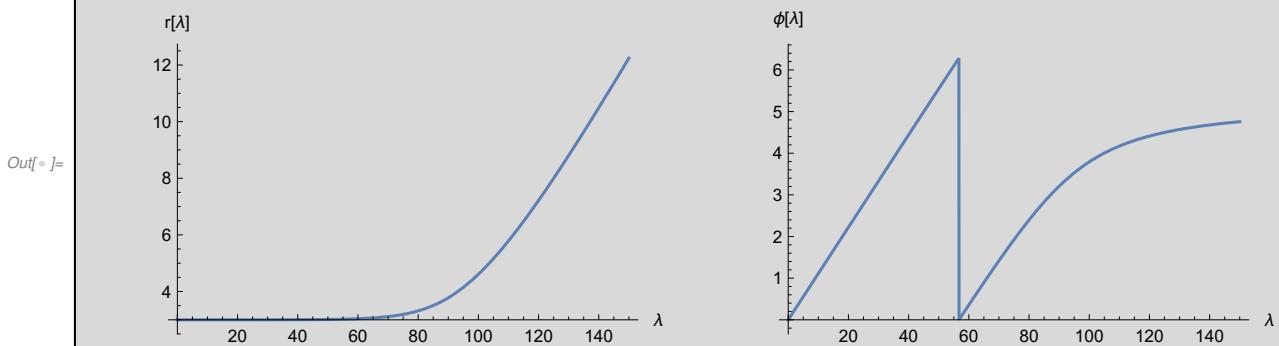
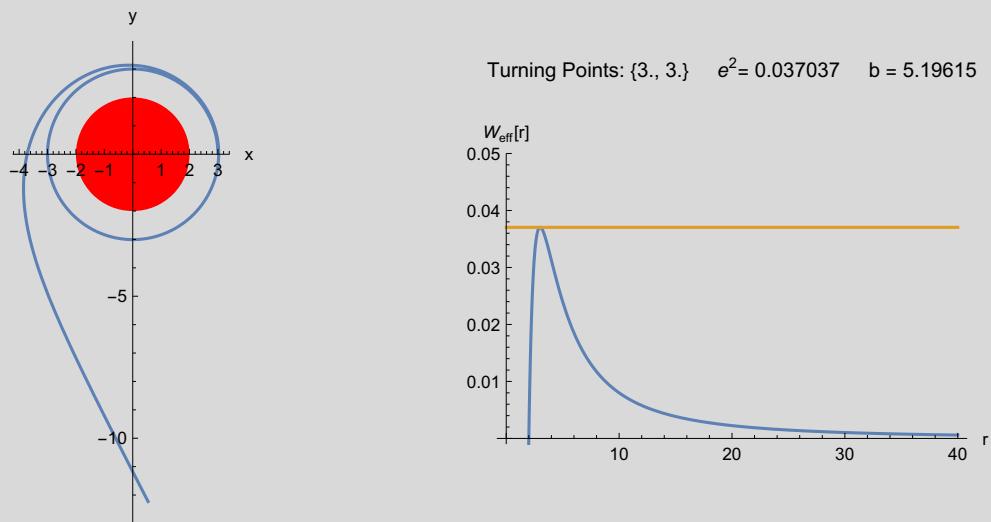


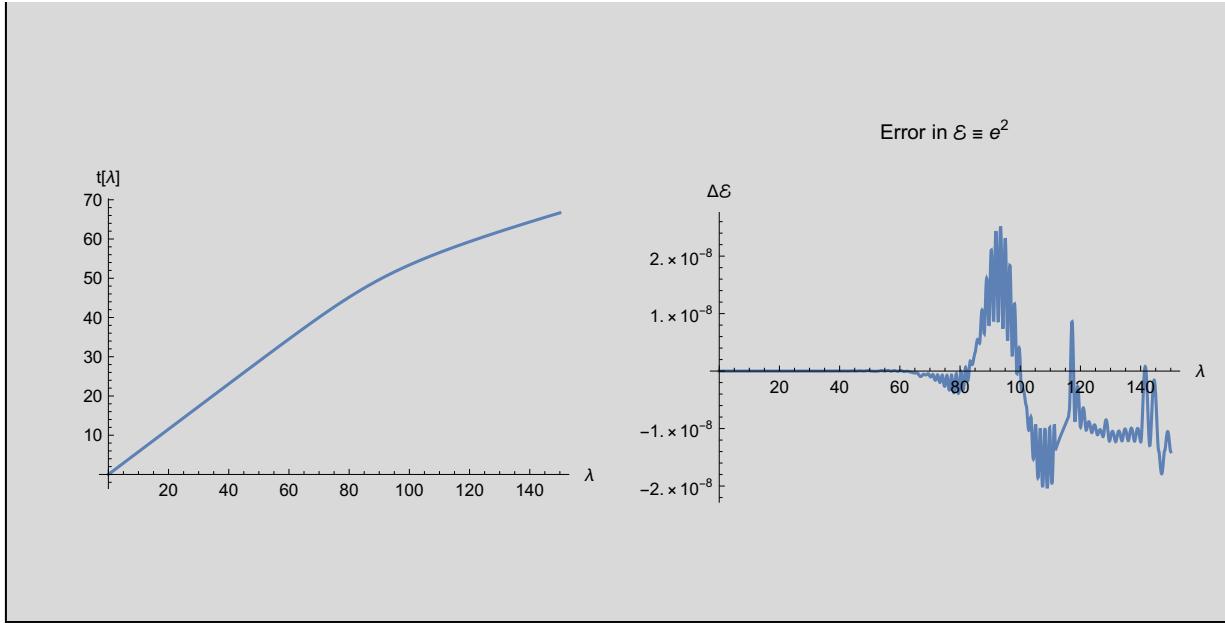
Instability of the circular orbits:

Slightly perturbed away from the  $r_{\max} = 3$ ,  $W_{\max} = \frac{1}{27}$  circular orbit

```
In[1]:= r1 = 3.0001 ; λmax = 150 ;
r0 = 3.0001 ; radialdirection = -1; (* Set it to +/- 1. The sign of ̇r(0) *)
Energy = Weff[r1]; e = √[Energy] ; b = 1/e;
v0 = radialdirection √[(Energy - Weff[r0])];
ϕ0 = 0;

sol = solveGeodesicEqs[e, ϕ0, r0, v0, λmax]; gg = plotResults[All]
```





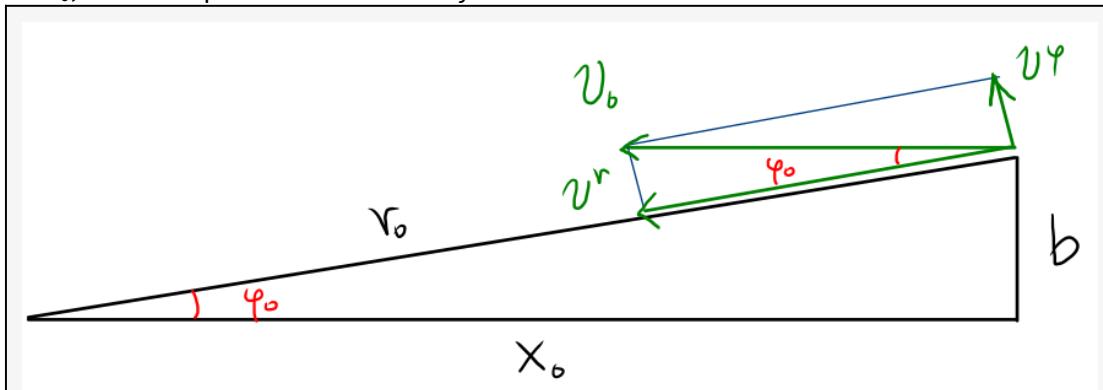
## Scattering - Deflection of light

Set:

$b$ : the impact parameter

$x_0$ : initial x coordinate

$v_0$ :  $v_0$ , the x-component of the velocity



Then:

$$r_0 = \sqrt{x_0^2 + b^2}$$

$$v_0 = 1 = \frac{dx}{dt} \quad (\text{far from } r=0)$$

$$v^r =$$

$$-v_0 \cos(\phi_0) = \frac{dr}{dt} = \frac{dr}{d\lambda} / \frac{dt}{d\lambda} = \dot{r}(0) / \left( e \left( 1 - \frac{2}{r} \right)^{-1} \right) \Rightarrow \dot{r}(0) = -e \left( 1 - \frac{2}{r} \right)^{-1} \cos(\phi_0) \Rightarrow \dot{r}(0) = -\frac{1}{b} \left( 1 - \frac{2}{r} \right)^{-1} \cos(\phi_0)$$

$$e^2 = \dot{r}_0^2 + W_{\text{eff}}(r_0)$$

To have finite deflection angle, we must have  $b > \sqrt{27} \approx 5.196$

```

b = 8; x0 = 500.; λmax = 12 000 ;

r0 = √(x02 + b2 );
ϕ0 = ArcTan[b/x0];
v0 = -1/b (1 - 2/r0)-1 Cos[ϕ0];

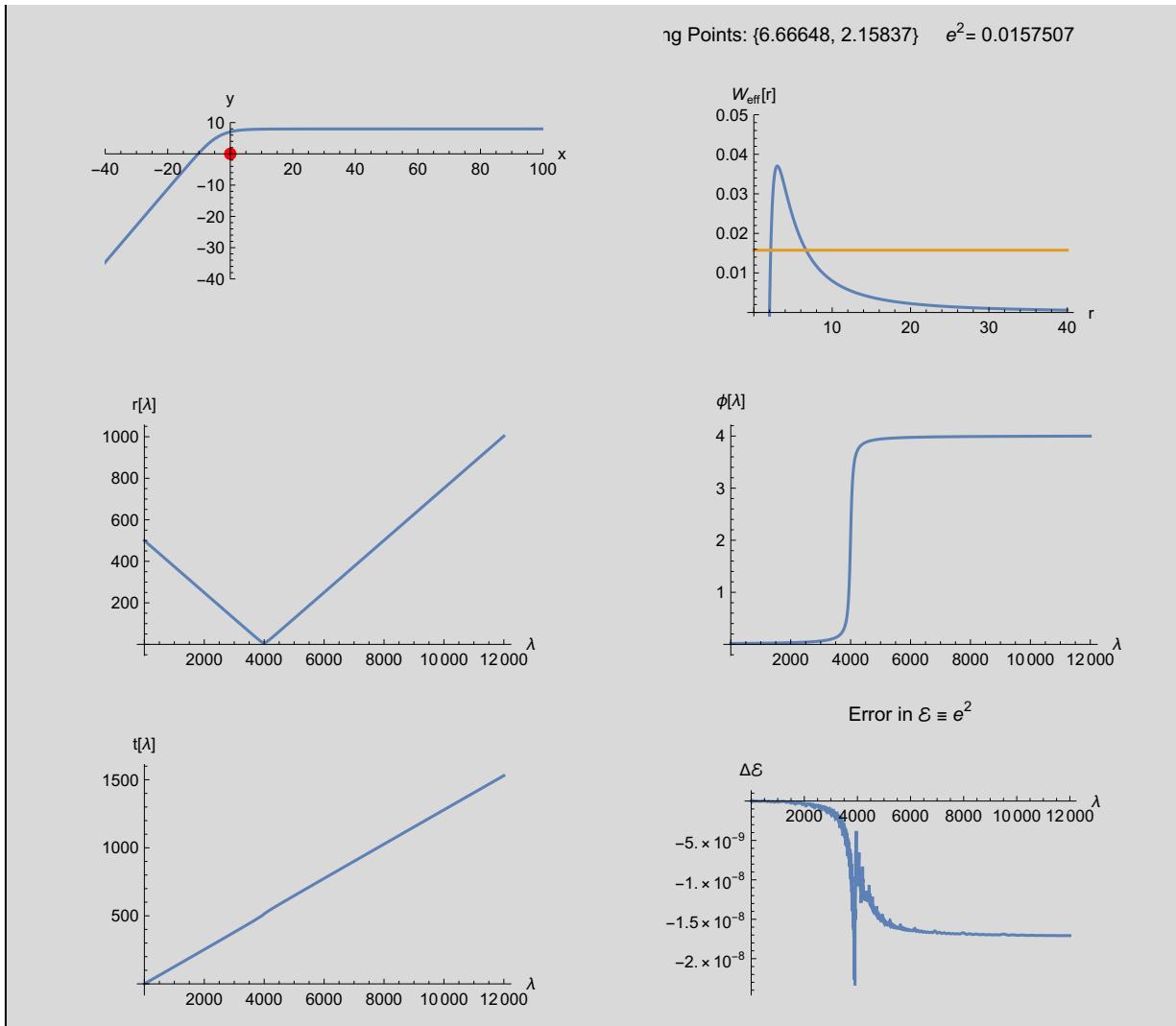
Energy = v02 + Weff[r0];
e = √Energy ; b = 1/e; (* b must be redefined,
it is only approximately equal to the impact parameter *)

sol = solveGeodesicEqs[e, ϕ0, r0, v0, λmax];
gg = plotResults[{{-40, 100}, {-40, 10}}];

Print[
  "Scattering Angle: (degrees)\n",
  "θ= tan-1vy(∞)/vx(∞) = ",
  Δθ = 360 + 180/π (ArcTan[r'[λmax] Cos[ϕ[λmax]] - ϕ'[λmax] r[λmax] Sin[ϕ[λmax]] /. sol,
    r'[λmax] Sin[ϕ[λmax]] + ϕ'[λmax] r[λmax] Cos[ϕ[λmax]] /. sol]) , "\n",
  "δθdeflection= ", Δθ - 180
];
Show[gg]

Scattering Angle: (degrees)
θ= tan-1vy(∞)/vx(∞) = {229.575}
δθdeflection= {49.5747}

```



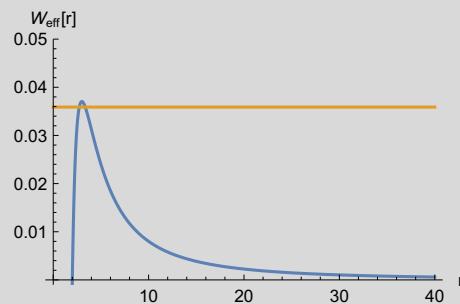
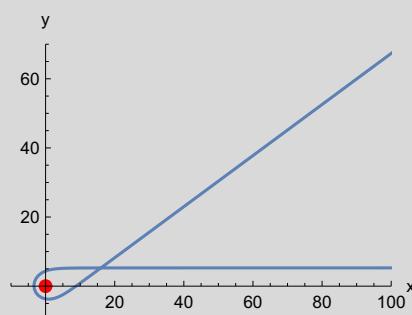
```
In[6]:= b = 5.3; x0 = 500.; λmax = 12000;
```

$$\begin{aligned} r_0 &= \sqrt{x_0^2 + b^2}; \\ \phi_0 &= \text{ArcTan}\left[\frac{b}{x_0}\right]; \\ v_0 &= -\frac{1}{b} \left(1 - \frac{2}{r_0}\right)^{-1} \cos[\phi_0]; \end{aligned}$$

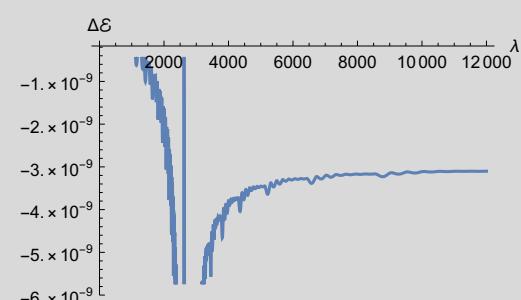
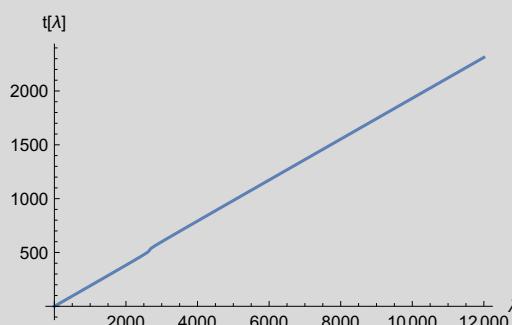
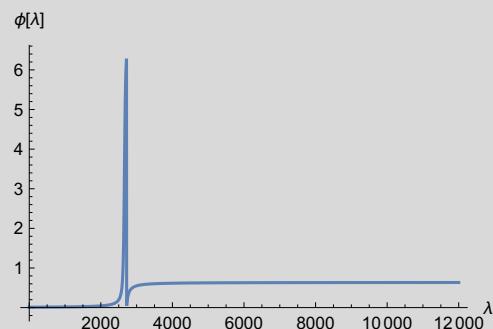
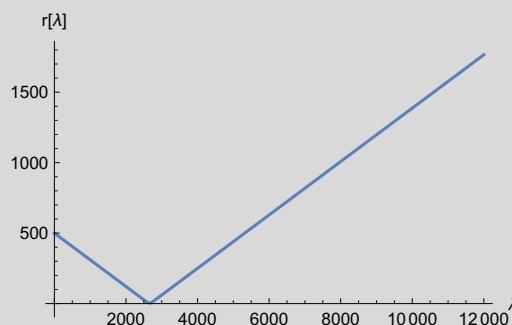
$$\text{Energy} = v_0^2 + W_{\text{eff}}[r_0]; e = \sqrt{\text{Energy}}; b = 1/e;$$

```
sol = solveGeodesicEqs[e, φ0, r0, v0, λmax];
gg = plotResults[{{-10, 100}, {-10, 70}}]
```

Irning Points: {3.35401, 2.7308}  $e^2 = 0.0358863$   $b =$



```
Out[6]:=
```

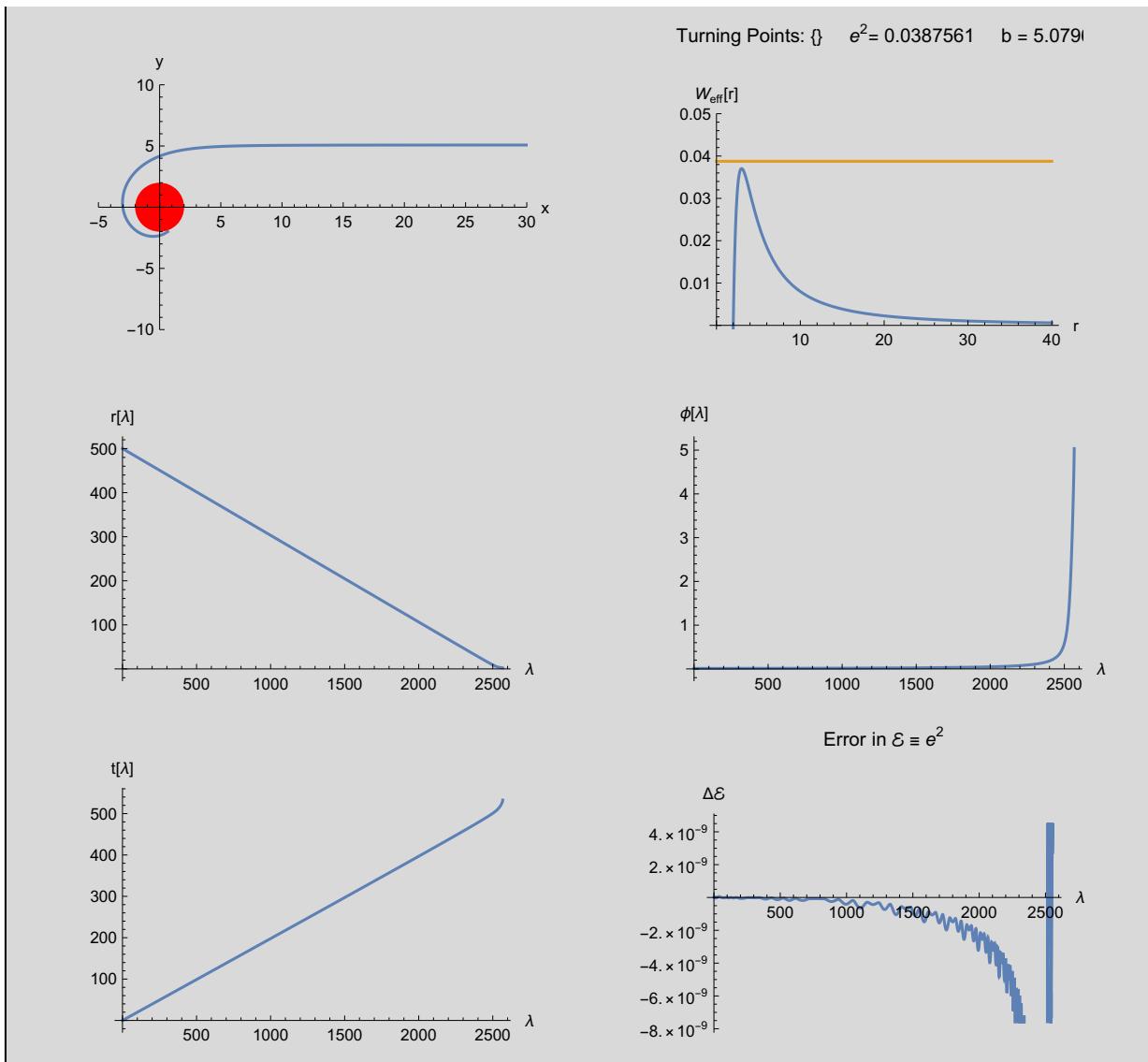


```
In[1]:= b = 5.1; x0 = 500.; λmax = 2567;

r0 = Sqrt[x0^2 + b^2];
ϕ0 = ArcTan[b/x0];
v0 = -1/b (1 - 2/r0)^(-1) Cos[ϕ0];

Weff[r_] := 1/r^2 (1 - 2/r);
Energy = v0^2 + Weff[r0];
e = Sqrt[Energy]; b = 1/e; (* b must be redefined,
it is only approximately equal to the impact parameter *)

sol = solveGeodesicEqs[e, ϕ0, r0, v0, λmax];
gg = plotResults[{{-5, 30}, {-10, 10}}]
```



```

In[1]:= b = 5.2170155293; x0 = 500.; λmax = 3000;

r0 = √(x0^2 + b^2);
ϕ0 = ArcTan[b/x0];
v0 = -1/b (1 - 2/r0)^(-1) Cos[ϕ0];

Energy = v0^2 + Weff[r0];
e = √Energy; b = 1/e; (* b must be redefined,
it is only approximately equal to the impact parameter *)

sol = solveGeodesicEqs[e, ϕ0, r0, v0, λmax];
gg = plotResults[{{-5, 30}, {-10, 10}}];

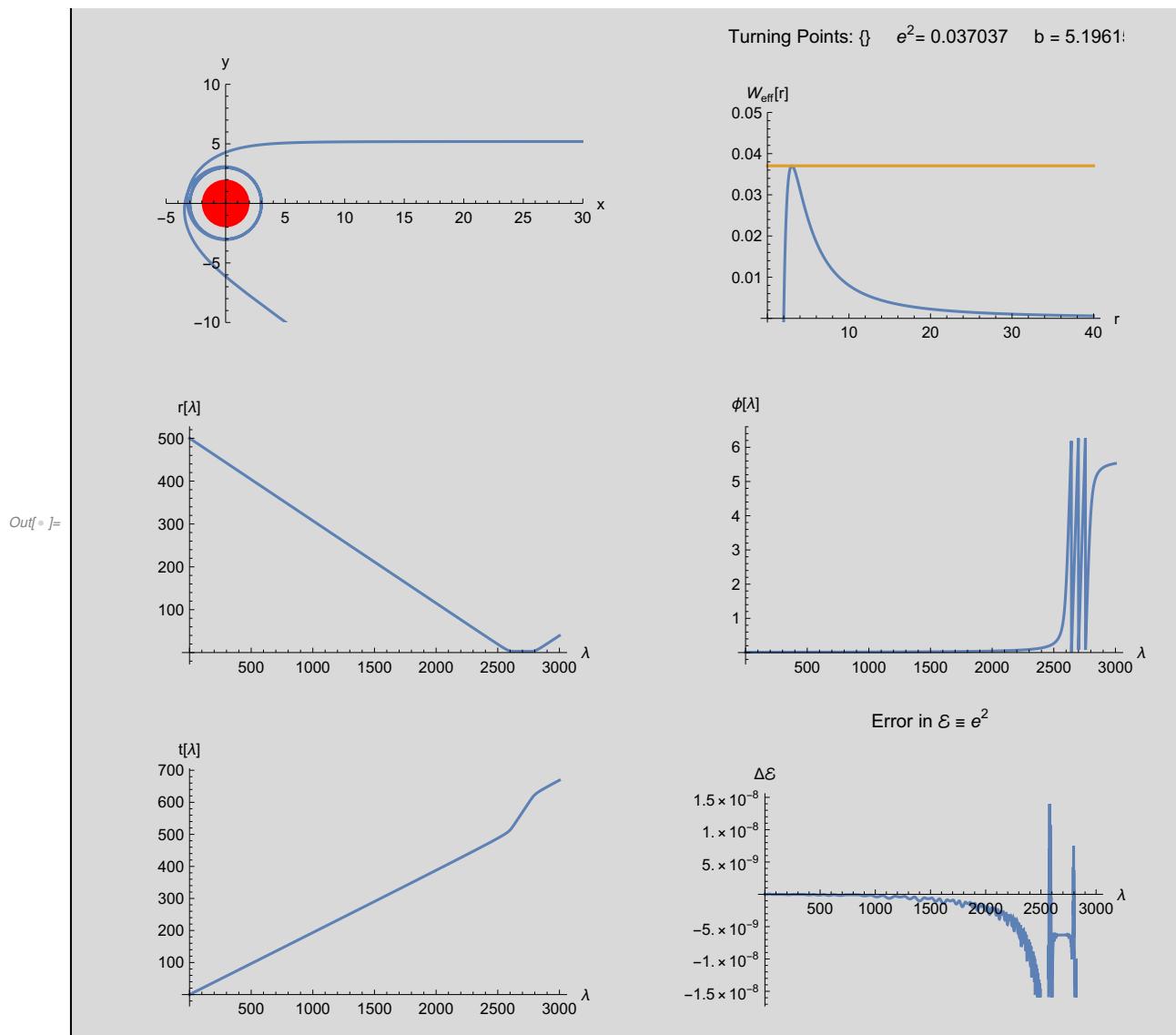
Print[
  "Scattering Angle: (degrees)\n",
  "θ= tan⁻¹(vy(∞)/vx(∞)) = ",
  Δθ = 360 + 180/π (ArcTan[r'[λmax] Cos[ϕ[λmax]] - ϕ'[λmax] r[λmax] Sin[ϕ[λmax]] /. sol,
    r'[λmax] Sin[ϕ[λmax]] + ϕ'[λmax] r[λmax] Cos[ϕ[λmax]] /. sol]), "\n",
  "δθdeflection= ", Δθ - 180
];
Show[gg]

```

Scattering Angle: (degrees)

$$\theta = \tan^{-1} \frac{v_y(\infty)}{v_x(\infty)} = \{324.304\}$$

$$\delta\theta_{\text{deflection}} = \{144.304\}$$



## Acknowledgements

This notebook has been programmed by Konstantinos Anagnostopoulos, Physics Department, National Technical University of Athens, Greece, while he was an instructor of the 4th year undergraduate course "General Relativity and Cosmology". It was created for fun, but it may turn out to be useful to everyone studying the General Theory of Relativity for the first time.

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